quantum!mechanics
mechanics!quantum
Schr“”odinger!equation
Hamiltonian
Schr“”odinger!equation!time independent
eigenfunction!energy

Chapter 32

Quantum mechanics, briefly

We start with a review of standard quantum mechanical concepts prerequisite to the derivation of the semiclassical trace formula.

In coordinate representation the time evolution of a quantum mechanical wave function is governed by the Schrödinger equation

$$i\hbar\frac{\partial}{\partial t}\psi(q,t) = \hat{H}(q,\frac{\hbar}{i}\frac{\partial}{\partial q})\psi(q,t), \tag{32.1}$$

where the Hamilton operator $\hat{H}(q,-i\hbar\partial_q)$ is obtained from the classical Hamiltonian by substitution $p \to -i\hbar\partial_q$. Most of the Hamiltonians we shall consider here are of form

$$H(q,p) = T(p) + V(q)\,, \qquad T(p) = p^2/2m\,, \tag{32.2}$$

describing dynamics of a particle in a D-dimensional potential $V(q)$. For time independent Hamiltonians we are interested in finding stationary solutions of the Schrödinger equation of the form

$$\psi_n(q,t) = e^{-iE_n t/\hbar}\phi_n(q), \tag{32.3}$$

where E_n are the eigenenergies of the time-independent Schrödinger equation

$$\hat{H}\phi(q) = E\phi(q)\,. \tag{32.4}$$

If the kinetic term can be separated out as in (32.2), the time-independent Schrödinger equation

$$-\frac{\hbar^2}{2m}\partial^2\phi(q) + V(q)\phi(q) = E\phi(q) \tag{32.5}$$

Hilbert!space

can be rewritten in terms of a local wavenumber

$$(\partial^2 + k^2(q))\phi = 0\,, \qquad \hbar^2 k(q) = \sqrt{2m(E - V(q))}\,. \tag{32.6}$$

For bound systems the spectrum is discrete and the eigenfunctions form an orthonormal,

$$\int dq\, \phi_n(q)\phi_m^*(q) = \delta_{nm}\,, \tag{32.7}$$

and complete,

$$\sum_n \phi_n(q)\phi_n^*(q') = \delta(q - q')\,, \tag{32.8}$$

set of functions in a Hilbert space. Here and throughout the text,

$$\int dq = \int dq_1 dq_2 ... dq_D. \tag{32.9}$$

For simplicity we will assume that the system is bound, although most of the results will be applicable to open systems, where one has complex resonances instead of real energies, and the spectrum has continuous components.

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A given wave function can be expanded in the energy eigenbasis

$$\psi(q,t) = \sum_n c_n e^{-iE_n t/\hbar}\phi_n(q)\,, \tag{32.10}$$

where the expansion coefficient c_n is given by the projection of the initial wave function $\psi(q,0)$ onto the nth eigenstate

$$c_n = \int dq\, \phi_n^*(q)\psi(q,0). \tag{32.11}$$

By substituting (32.11) into (32.10), we can cast the evolution of a wave function into a multiplicative form

$$\psi(q,t) = \int dq' K(q,q',t)\psi(q',0)\,,$$

with the kernel

$$K(q,q',t) = \sum_n \phi_n(q)\, e^{-iE_n t/\hbar}\phi_n^*(q') \tag{32.12}$$

called the quantum evolution operator, or the *propagator*. Applied twice, first for time t_1 and then for time t_2, it propagates the initial wave function from q' to q'', and then from q'' to q

$$K(q,q',t_1+t_2) = \int dq''\, K(q,q'',t_2)K(q'',q',t_1) \tag{32.13}$$

forward in time, hence the name "propagator". In non-relativistic quantum mechanics the range of q'' is infinite, meaning that the wave can propagate at any speed; in relativistic quantum mechanics this is rectified by restricting the propagation to the forward light cone.

Since the propagator is a linear combination of the eigenfunctions of the Schrödinger equation, it also satisfies the Schrödinger equation

$$i\hbar\frac{\partial}{\partial t}K(q,q',t) = \hat{H}(q,\frac{i}{\hbar}\frac{\partial}{\partial q})K(q,q',t)\,, \tag{32.14}$$

and is thus a wave function defined for $t \geq 0$; from the completeness relation (32.8) we obtain the boundary condition at $t = 0$:

$$\lim_{t\to 0_+} K(q,q',t) = \delta(q-q')\,. \tag{32.15}$$

The propagator thus represents the time evolution of a wave packet which starts out as a configuration space delta-function localized in the point q' at the initial time $t = 0$.

For time independent Hamiltonians the time dependence of the wave functions is known as soon as the eigenenergies E_n and eigenfunctions ϕ_n have been determined. With time dependence rendered "trivial", it makes sense to focus on the *Green's function*, the Laplace transformation of the propagator

$$G(q,q',E+i\epsilon) = \frac{1}{i\hbar}\int_0^\infty dt\, e^{\frac{i}{\hbar}Et-\frac{\epsilon}{\hbar}t}K(q,q',t) = \sum_n \frac{\phi_n(q)\phi_n^*(q')}{E-E_n+i\epsilon}\,. \tag{32.16}$$

Here ϵ is a small positive number, ensuring the existence of the integral. The eigenenergies show up as poles in the Green's function with residues corresponding to the wave function amplitudes. If one is only interested in the spectrum, one may restrict the considerations to the (formal) trace of the Green's function,

$$\mathrm{tr}\, G(q,q',E) = \int dq\, G(q,q,E) = \sum_n \frac{1}{E-E_n}\,, \tag{32.17}$$

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quantum!evolution
quantum!propagator
propagator
delta function!Dirac
Green's function!energy dependent
Laplace!transform
Green's function!trace

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IOP usage is: Green function
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32.2
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density of states!quantum
Green's function!density of states
density of states!Green's function

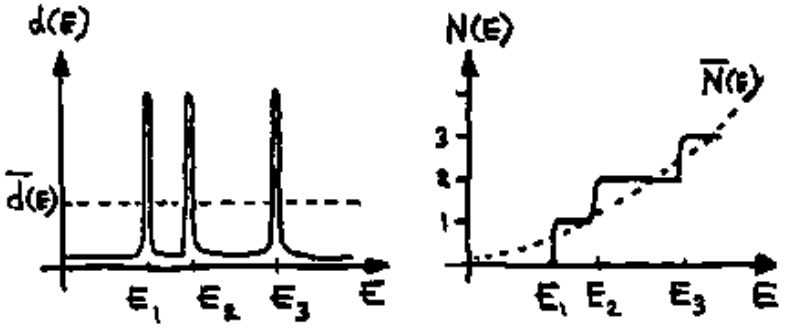

Figure 32.1: Schematic picture of **a)** the density of states $d(E)$, and **b)** the spectral staircase function $N(E)$. The dashed lines denote the mean density of states $\bar{d}(E)$ and the average number of states $\bar{N}(E)$ discussed in more detail in sect. 36.1.1.

where E is complex, with a positive imaginary part, and we have used the eigenfunction orthonormality (32.7). This trace is formal, since as it stands, the sum in (32.17) is often divergent. [1] We shall return to this point in sects. 36.1.1 and 36.1.2. [2]

A useful characterization of the set of eigenvalues is given in terms of the *density of states*, with a delta function peak at each eigenenergy, figure 32.1 (a),

$$d(E) = \sum_n \delta(E - E_n). \tag{32.18}$$

32.1
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Using the identity

$$\delta(E - E_n) = -\lim_{\epsilon \to +0} \frac{1}{\pi} \mathrm{Im} \frac{1}{E - E_n + i\epsilon} \tag{32.19}$$

we can express the density of states in terms of the trace of the Green's function, that is

$$d(E) = \sum_n \delta(E - E_n) = -\lim_{\epsilon \to 0} \frac{1}{\pi} \mathrm{Im}\, \mathrm{tr}\, G(q, q', E + i\epsilon). \tag{32.20}$$

☞ sect. 36.1.1

As we shall see after "some" work, a semiclassical formula for right hand side of this relation will yield the quantum spectrum in terms of periodic orbits.

[1] Whelan: The convergence depends on the density of states. For a 1D infinite square well the energies are proportional to n^2 and the equation above would converge. I think (but may be wrong, I am going on memory) that there is a complementarity. If the trace of the propagator converges the trace of the Green's function does not and vice-versa. By the way, you may want to write down the trace of the propagator as $K(t) = \sum_{n} exp(-iE_n t/\hbar)$.

[2] Tanner: Cite Voros, Sieber ...

The density of states can be written as the derivative $d(E) = \mathrm{d}N(E)/\mathrm{d}E$ of the *spectral staircase* function

spectral!staircase
staircase!spectral
Heaviside function

$$N(E) = \sum_n \Theta(E - E_n) \tag{32.21}$$

which counts the number of eigenenergies below E, figure 32.1 (b). Here Θ is the Heaviside function

$$\Theta(x) = 1 \quad \text{if } x > 0; \quad \Theta(x) = 0 \quad \text{if } x < 0\,. \tag{32.22}$$

add spect determinant here

The spectral staircase is a useful quantity in many contexts, both experimental and theoretical. This completes our lightning review of quantum mechanics.

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