Chapter 35

Noise

> He who establishes his argument by noise and command shows that his reason is weak.
> M. de Montaigne

(G. Vattay and P. Cvitanović)

This chapter (which reader can safely skip on the first reading) is about noise, how it affects classical dynamics, and the ways it mimicks quantum dynamics.

fast track:
chapter 36, p. 663

Why - in a study of deterministic and quantum chaos - start discussing noise? First, in physical settings any dynamics takes place against a noisy background, and whatever prediction we might have, we have to check its robustness to noise. Second, as we show in this chapter, to the leading order in noise strength the semiclassical Hamilton-Jacobi formalism carries over to weakly stochastic flows in totto. As classical noisy dynamics is more intuitive than quantum dynamics, this exercise helps demistify some of the formal machinery of semiclassical quantization. Surprisingly, symplectic structure emerges here not as a deep principle of mechanics, but an artifact of the leading approximation to quantum/noisy dynamics, not respected by higher order corrections. The same is true of semiclasical quantum dynamics; higher corrections do not respect canonical invariance. Third, the variational principle derived here will be refashioned into a powerful tool for determining periodic orbits in chapter 37.

We start by deriving the continuity equation for purely deterministic, noiseless flow, and then incorporate noise in stages: diffusion equation, Langevin equation, Fokker-Planck equation, Hamilton-Jacobi formulation, stochastic path integrals.

material invariant
passive scalar
conservation!equation
density
continuity equation

35.1 Deterministic transport

(E.A. Spiegel and P. Cvitanović)

Fluid dynamics is about physical flows of media with continuous densities. On the other hand, the flows in state spaces of dynamical systems frequently require more abstract tools. To sharpen our intuition about those, it is helpful to outline the more tangible fluid dynamical vision.

Consider first the simplest property of a fluid flow called *material invariant*. A material invariant $I(x)$ is a property attached to each point x that is preserved by the flow, $I(x) = I(f^t(x))$; for example, at this point a green particle (more formally: a *passive scalar*) is embedded into the fluid. As $I(x)$ is invariant, its total time derivative vanishes, $\dot{I}(x) = 0$. Written in terms of partial derivatives this is the *conservation equation* for the material invariant

$$\partial_t I + v \cdot \partial I = 0\,. \tag{35.1}$$

Let the *density* of representative points be $\rho(x,t)$. The manner in which the flow redistributes $I(x)$ is governed by a partial differential equation whose form is relatively simple because the representative points are neither created nor destroyed. This conservation property is expressed in the integral statement

$$\partial_t \int_V dx\, \rho I = -\int_{\partial V} d\sigma\, \hat{n}_i v_i \rho I\,,$$

where V is an arbitrary volume in the state space $\mathcal{M}$, ∂V is its surface, $\hat{n}$ is its outward normal, and repeated indices are summed over throughout. The divergence theorem turns the surface integral into a volume integral,

$$\int_V [\partial_t(\rho I) + \partial_i(v_i \rho I)]\, dx = 0\,,$$

where ∂_i is the partial derivative operator with respect to x_i. Since the integration is over an arbitrary volume, we conclude that

$$\partial_t(\rho I) + \partial_i(\rho I v_i) = 0\,. \tag{35.2}$$

The choice $I \equiv 1$ yields the *continuity equation* for the density:

$$\partial_t \rho + \partial_i(\rho v_i) = 0\,. \tag{35.3}$$

We have used here the language of fluid mechanics to ease the visualization, but, as we already saw in the discussion of infinitesimal action of the Perron-Frobenius operator (9.24), continuity equation applies to any deterministic state space flow.

35.2 Brownian difussion

Fick law
diffusion!equation
Dirac delta function
Einstein!diffusion formula

Consider tracer molecules, let us say green molecules, embedded in a denser gas of transparent molecules. Assume that the density of tracer molecules ρ compared to the background gas density is low, so we can neglect green-green collisions. Each green molecule, jostled by frequent collisions with the background gas, executes its own Brownian motion. The molecules are neither created nor destroyed, so their number within an arbitrary volume V changes with time only by the current density j_i flow through its surface ∂V (with $\hat{n}$ its outward normal):

$$\partial_t \int_V dx\, \rho = -\int_{\partial V} d\sigma\, \hat{n}_i j_i \,. \tag{35.4}$$

The divergence theorem turns this into the conservation law for tracer density:

$$\partial_t \rho + \partial_i j_i = 0 \,. \tag{35.5}$$

The tracer density ρ is defined as the average density of a "material particle", averaged over a subvolume large enough to contain many green (and still many more background) molecules, but small compared to the macroscopic observational scales. What is j? If the density is constant, on the average as many molecules leave the material particle volume as they enter it, so a reasonable phenomenological assumption is that the *average* current density (*not* the individual particle current density ρv_i in (35.3)) is driven by the density gradient

$$j_i = -D\frac{\partial \rho}{\partial x_i} \,. \tag{35.6}$$

This is the *Fick law*, with the diffusion constant D a phenomenological parameter. For simplicity here we assume that D is a scalar; in general $D \to D_{ij}(x,t)$ is a space- and time-dependent tensor. Substituting this j into (35.5) yields the *diffusion equation*

$$\frac{\partial}{\partial t}\rho(x,t) = D\frac{\partial^2}{\partial x^2}\rho(x,t) \,. \tag{35.7}$$

This linear equation has an exact solution in terms of an initial Dirac delta density distribution, $\rho(x,0) = \delta(x - x_0)$,

$$\rho(x,t) = \frac{1}{(4\pi D t)^{d/2}} e^{-\frac{(x-x_0)^2}{4Dt}} = \frac{1}{(4\pi D t)^{d/2}} e^{-\frac{\dot{x}^2}{4D}t} \,, \tag{35.8}$$

check factor of 4 here

noise!Gaussian
Gaussian!probability density
probability density!Gaussian
white noise
noise!white
Langevin equation

in the spirit of the quantum free particle propagation of sect. 34.2.2. The average distance covered in time t obeys the Einstein diffusion formula

change this - Ein is specific to Brow diffusion

$$\left\langle (x-x_0)^2 \right\rangle_t = \int dx\, \rho(x,t)(x-x_0)^2 = 2Dt\,. \tag{35.9}$$

35.3 Weak noise

> The connection between path integration and Brownian motion is so close that they are nearly indistingushable. Unfortunately though, like a body and its mirror image, the sum over paths for Brownian motion is a theory having substance, while its path integral image exists mainly in the eye of the beholder.
>
> L. S. Schulman

So far we have considered tracer molecule dynamics which is purely Brownian, with no deterministic "drift". Consider next a deterministic flow $\dot{x} = v(x)$ perturbed by a stochastic term $\xi(t)$,

$$\dot{x} = v(x) + \xi(t)\,. \tag{35.10}$$

check factor of 4

Assume that $\xi(t)$'s fluctuate around $[\dot{x} - v(x)]$ with a Gaussian probability density

$$p(\xi, \delta t) = \left(\frac{\delta t}{4\pi D}\right)^{d/2} e^{-\frac{\xi^2}{4D}\delta t}\,, \tag{35.11}$$

explain what average?

and are uncorrelated in time (white noise)

$$\left\langle \xi(t)\xi(t') \right\rangle = 2dD\delta(t-t')\,. \tag{35.12}$$

The normalization factors in (35.8) and (35.11) differ, as $p(\xi, \delta t)$ is a probability density for velocity ξ, and $\rho(x,t)$ is a probability density for position x. The material particle now drifts along the trajectory $x(t)$, so the velocity diffusion follows (35.8) for infinitesimal time δt only. As $D \to 0$, the distribution tends to the (noiseless, deterministic) Dirac delta function.

An example is the Langevin equation for a Brownian particle, in which one replaces the Newton's equation for force by two counter-balancing forces: random accelerations $\xi(t)$ which tend to smear out a particle trajectory, and a damping term which drives the velocity to zero.

explain what $|v(x)| > |\xi|$ means

The phenomenological Fick law current (35.6) is now a sum of two components, the material particle center-of-mass deterministic drift $v(x)$ and the weak noise term

$$j_i = \rho v_i - D\frac{\partial \rho}{\partial x_i}\,, \tag{35.13}$$

Substituting this j into (35.5) yields the *Fokker-Planck equation*

Fokker-Planck equation
continuity equation

$$\partial_t \rho + \partial_i(\rho v_i) = D\,\partial^2 \rho. \tag{35.14}$$

The left hand side, $d\rho/dt = \partial_t \rho + \partial \cdot (\rho v)$, is deterministic, with the continuity equation (35.3) recovered in the weak noise limit $D \to 0$. The right hand side describes the diffusive transport in or out of the material particle volume. If the density is lower than in the immediate neighborhood, the local curvature is positive, $\partial^2 \rho > 0$, and the density grows. Conversely, for negative curvature diffusion lowers the local density, thus smoothing the variability of ρ. Where is the density going globally?

If the system is bound, the probability density vanishes sufficiently fast outside the central region, $\rho(x,t) \to 0$ as $|x| \to \infty$, and the total probability is conserved

$$\int dx\, \rho(x,t) = 1\,.$$

Any initial density $\rho(x,0)$ is smoothed by diffusion and with time tends to the invariant density

$$\rho_0(x) = \lim_{t\to\infty} \rho(x,t)\,, \tag{35.15}$$

an eigenfunction $\rho(x,t) = e^{st}\,\rho_0(x)$ of the time-independent Fokker-Planck equation

$$\left(\partial_i v_i - D\,\partial^2 + s_\alpha\right)\rho_\alpha = 0\,, \tag{35.16}$$

with vanishing eigenvalue $s_0 = 0$. Provided the noiseless classical flow is hyperbolic, in the vanishing noise limit the leading eigenfunction of the Fokker-Planck equation tends to natural measure (9.16) of the corresponding deterministic flow, the leading eigenvector of the Perron-Frobenius operator.

If the system is open, there is a continuous outflow of probability from the region under study, the leading eigenvalue is contracting, $s_0 < 0$, and the density of the system tends to zero. In this case the leading eigenvalue s_0 of the time-independent Fokker-Planck equation (35.16) can be interpreted by saying that a finite density can be maintained by pumping back probability into the system at a constant rate $\gamma = -s_0$. The value of γ for which any initial probability density converges to a finite equilibrium density is called the *escape rate*. In the noiseless limit this coincides with the deterministic escape rate (10.15).

We have introduced noise phenomenologically, and used the weak noise assumption in retaining only the first derivative of ρ in formulating the Fick law (35.6) and including noise additively in (35.13). A full theory of stochastic ODEs is much subtler, but this will do for our purposes.

Hamilton!principal function
matrix!stability
stability!matrix

35.4 Weak noise approximation

In the spirit of the WKB approximation, we shall now study the evolution of the probability distribution by rewriting it as

$$\rho(x,t) = e^{\frac{1}{2D}R(x,t)}\,. \tag{35.17}$$

The time evolution of R is given by

$$\partial_t R + v\partial R + (\partial R)^2 = D\partial v + D\partial^2 R\,.$$

Consider now the weak noise limit and drop the terms proportional to D. The remaining equation

$$\partial_t R + H(x,\partial R) = 0$$

is the Hamilton-Jacobi equation (34.5). The function R can be interpreted as the Hamilton's principal function, corresponding to the Hamiltonian

$$H(x,p) = p\,v(x) + p^2/2\,,$$

with the Hamilton's equations of motion

$$\begin{aligned} \dot{x} &= \partial_p H = v + p \\ \dot{p} &= -\partial_x H = -A^T p\,, \end{aligned} \tag{35.18}$$

where A is the stability matrix (4.3)

$$A_{ij}(x) = \frac{\partial v_i(x)}{\partial x_j}\,.$$

The noise Lagrangian (34.13) is then

$$L(x,\dot{x}) = \dot{x}\cdot p - H = \frac{1}{2}\left[\dot{x} - v(x)\right]^2\,. \tag{35.19}$$

We have come the full circle - the Lagrangian is the exponent of our assumed Gaussian distribution (35.11) for noise $\xi^2 = [\dot{x} - v(x)]^2$. What is the meaning of this Hamiltonian, Lagrangian? Consider two points x_0 and x. Which noisy path is the most probable path that connects them in time t? The probability of a given path $\mathcal{P}$ is given by the probability of the noise sequence $\xi(t)$ which generates the path. This probability is proportional to the product of the noise probability functions (35.11) along the path, and

eck factors of 2

Wiener integral
variational principle

the total probability for reaching x from x_0 in time t is given by the sum over all paths, or the stochastic path integral (Wiener integral)

$$\begin{aligned} P(x,x_0,t) &\sim \sum_{\mathcal{P}} \prod_j p(\xi(\tau_j),\delta\tau_j) = \int \prod_j d\xi_j \left(\frac{\delta\tau_j}{2\pi D}\right)^{d/2} e^{-\frac{\xi(\tau_j)^2}{2D}\delta\tau_i} \\ &\to \frac{1}{Z}\sum_{\mathcal{P}} \exp\left(-\frac{1}{2D}\int_0^t d\tau\, \xi^2(\tau)\right) , \end{aligned} \tag{35.20}$$

where $\delta\tau_i = \tau_i - \tau_i$, and the normalization constant is

$$\frac{1}{Z} = \lim \prod_i \left(\frac{\delta\tau_i}{2\pi D}\right)^{d/2} .$$

The most probable path is the one maximizing the integral inside the exponential. If we express the noise (35.10) as

$$\xi(t) = \dot{x}(t) - v(x(t)) ,$$

the probability is maximized by the variational principle

$$\min \int_0^t d\tau [\dot{x}(\tau) - v(x(\tau))]^2 = \min \int_0^t L(x(\tau),\dot{x}(\tau)) d\tau .$$

By the standard arguments, for a given x, x' and t the the probability is maximized by a solution of Hamilton's equations (35.18) that connects the two points $x_0 \to x'$ in time t.

⇓PRIVATE

35.5 Straying from the True Path: the cost

Penalizing deviations from the True Path?

Newton descent derived from path integral for a weakly stochastic Langevin flow

stepwise error

$$F_i = \tilde{x}_{i+1} - f^{\Delta t_i}(\tilde{x}_i)$$

where $\tilde{x}_{i+1}$ is the next path point, and $f^{\Delta t_i}(\tilde{x}_i)$ is the point to which the flow transports $\tilde{x}_i$.

Cost function

$$F^2 = \sum_{i=0}^{N} \frac{1}{\Delta t_i} \left(\tilde{x}_{i+1} - f^{\Delta t_i}(\tilde{x}_i)\right)^2$$

Dimensionally $F^2 = [\text{lenght}]^2/[\text{time}] = \textit{diffusion}$

rewrite as

$$\begin{aligned} F^2 &= \sum_{i=0}^{N} \frac{1}{\Delta t_i} \left(\tilde{x}_{i+1} - \tilde{x}_i - f^{\Delta t_i}(\tilde{x}_i) - \tilde{x}_i\right)^2 \\ &= \sum_{i=0}^{N} \Delta t_i \left(\frac{\tilde{x}_{i+1} - \tilde{x}_i}{\Delta t_i} - \frac{f^{\Delta t_i}(\tilde{x}_i) - \tilde{x}_i}{\Delta t_i}\right)^2 \end{aligned}$$

$\Delta t_i \to 0$ limit:

$$F^2 = \int_0^T dt \left(\frac{d\tilde{x}}{dt} - v(\tilde{x})\right)^2$$

where T = cycle period

$$F^2 = \int_0^T dt \left(\frac{d\tilde{x}}{dt} - v(\tilde{x})\right)^2$$

$$= = \int_0^T dt \left(\frac{ds}{dt}\right)^2 \left(\tilde{v}(\tilde{x}) - v(\tilde{x})\frac{dt}{ds}\right)^2, \quad (35.21)$$

where T = cycle period

to relate time t and the path parameter s consider the euclidean lengths

True Path: $dx = vdt$
Strayed: $d\tilde{x} = \tilde{v}ds$

related by projecting path step $\tilde{v}ds$ onto the true trajectory step vdt

$$\frac{|v|dt}{|\tilde{v}|ds} = \cos\theta\,, \quad (35.22)$$

where θ is the angle between v to $\tilde{v}$

$$\cos\theta = \frac{\tilde{v}\cdot v}{|\tilde{v}||v|} \quad (35.23)$$

and thus

$$\frac{dt}{ds} = \frac{\tilde{v}\cdot v}{v^2}\,. \quad (35.24)$$

$$F^2 = \int_{t_i=0}^{T} dt \left(\frac{v^2}{\tilde{v}\cdot v}\right)^2 \left(\tilde{v}(\tilde{x}) - v(\tilde{x})\frac{\tilde{v}\cdot v}{v^2}\right)^2 . \tag{35.25}$$

Langevin equation
Onsager-Machlup
van Kampen, N. G.

decomposing

$$\begin{aligned} \tilde{v} &= (1 - \frac{v\otimes v}{v^2} + \frac{v\otimes v}{v^2})\tilde{v} \\ &= = P_\perp v + v\frac{v\cdot\tilde{v}}{v^2} \end{aligned} \tag{35.26}$$

where

$$P_\perp = 1 - \frac{v\otimes v}{v^2}$$

projects any vector components transverse to v

Parametrized by the Euclidean length s measured along the guess path L the cost function is

$$F^2 = \int_{s_i}^{s_f} ds\, \frac{v^2}{v\cdot\tilde{v}}\,(P_\perp\tilde{v})^2 \tag{35.27}$$

The projection operator $P_\perp$ in the cost function penalizes the local misorientation of the path by projecting on the directions perpendicular to the dynamical flow vector.

⇑PRIVATE

Commentary

Remark 35.1 Literature. The theory of stochastic processes is a vast subject, spanning over centuries and over disciplines ranging from pure mathematics to impure finance. We enjoyed reading van Kampen classic [35.1], especially his railings against those who blunder carelessly into nonlinear landscapes (with this chapter we join the list of van Kampen's sinners). A more specialized monograph like Risken's [35.2] will do just as well. The "Langevin equation" introduces noise and damping only into the acceleration of Newton's equations; here we are considering more general stochastic differential equations in the weak noise limit. Onsager-Machlup seminal paper [35.18] was the first to introduce a variational method - the "principle of least dissipation" - based on the Lagrangian of form (35.19). This paper deals only with a finite set of linearly damped thermodynamic variables. Here the setting is much more general: we study fluctuations over a state space varying velocity field $v(x)$. Schulman's monograph [35.11] contains a very readable summary of Kac's [35.12] exposition of Wiener's integral over stochastic paths.

⇓PRELIMINARY

Remark 35.2 Operator ordering. The Itó and Stratanovich operator ordering issues arise in the order beyond the leading approximation considered here. [1]

not true? rethink

⇑PRELIMINARY

[1] Mason: make sure no ref. refer to volume 1

noise!Gaussian

Résumé

When a deterministic trajectory is smeared out under the influence of Gaussian noise of strength D, the deterministic dynamics is recovered in the weak noise limit $D \to 0$. The effect of the noise can be taken into account by adding noise corrections to the classical trace formula. [2]

Add Foged by reference

recheck whether something can be used from chapter/refsQmnoise

References

[35.1] N. G. van Kampen, *Stochastic Processes in Physics and Chemistry* (North Holland, Amsterdam, 1981).

[35.2] H. Risken, *The Fokker-Planck equation: Methods of solution and applications* (Springer-Verlag, New York, 1989).

[35.3] W. Dittrich and M. Reuter, *Classical and Quantum Dynamics: From Classical Paths to Path Integrals* (Springer-Verlag, Berlin 1994).

[35.4] E. M. Lifshitz and L. P. Pitaevskii, *Physical Kinetics* (Pergamon, London 1981).

[35.5] C. Itzykson and J.-M. Drouffe, *Statistical field theory* (Cambridge U. Press, 1991).

[35.6] V. Ambegaokar, *Reasoning about luck; probability and its uses in physics* (Cambridge Univ. Press, Cambridge 1996).

[35.7] B. Sakita, *Quantum theory of many variable systems and fields* (World Scientific, Singapore 1985).

[35.8] G. E. Uhlenbeck, G. W. Ford and E. W. Montroll, *Lectures in Statistical Mechanics* (Amer. Math. Soc., Providence R.I., 1963).

[35.9] M. Kac, "Random walk and the theory of Brownian motion", (1946), reprinted in ref. [35.10].

[35.10] N. Wax, ed., *Selected Papers on Noise and Stochastic Processes* (Dover, New York 1954).

[35.11] L. S. Schulman, *Techniques and Applications of Path Integration* (Wiley-Interscience, New York 1981).

[35.12] M. Kac, *Probability and Related Topics in Physical Sciences* (Wiley-Interscience, New York 1959).

[35.13] E. Nelson, *Quantum Fluctuations* (Princeton Univ. Press 1985).

[35.14] H. Kunita, *Stochastic Flows and Stochastic Differential Equations* (Cambridge Univ. Press, 1990).

[35.15] H. Haken and G. Mayer-Kress, *Z. f. Physik* **B 43**, 185 (1981).

[35.16] M. Roncadelli, *Phys. Rev.* **E 52**, 4661 (1995).

[35.17] G. Ryskin, *Phys. Rev.* **E 56**, 5123 (1997).

ref. [35.19]

[35.18] L. Onsager and S. Machlup, *Phys. Rev.* **91**, 1505, 1512 (1953).

[35.19] Lord Rayleigh, *Phil. Mag.* **26**, 776 (1913).

[35.20] E. Gozzi, M. Reuter and W. D. Thacker, *Phys. Rev.* **D 40**, 3363 (1989).

[35.21] E. Gozzi and M. Reuter, *Phys. Lett.* **233B**, 383 (1989); **238B**, 451 (1990); **240B**, 137 (1990).

[35.22] E. Gozzi, M. Reuter and W. D. Thacker, *Phys. Rev.* **D 46**, 757 (1992).

[35.23] E. Gozzi, M. Reuter and W. D. Thacker, *Chaos, Solitons and Fractals* **2**, 441 (1992).

[35.24] Benzi et al., *Jour. Phys.* **A 18**, 2157 (1985).

[35.25] R. Graham, *Europhys. Lett.* **5**, 101 (1988).

[35.26] R. Graham and T. Tél, *Phys. Rev. A* **35**, 1382 (1987).

[35.27] R. Benzi, G. Paladin, G. Parisi and A. Vulpiani, *J. Phys.* **A 18**, 2157 (1985).

[35.28] L. Arnold and V. Wihstutz, *Lyapunov exponents*, *Lecture Notes in Math.* **1186** (Springer-Verlag, New York 1986).

[35.29] M. V. Feigelman and A. M. Tsvelik, *Sov. Phys. JETP* **56**, 823 (1982).

[35.30] G. Parisi and N. Sourlas, *Nucl. Phys.* **B206**, 321 (1982).

[35.31] F. Langouche et al., *Functional integration and semiclassical expansion* (Reidel, Dordrecht 1982).

[35.32] O. Cepas and J. Kurchan, "Canonically invariant formulation of Langevin and Fokker-Planck equations", *European Phys. J.* **B 2**, 221 (1998), `cond-mat/9706296`.

[35.33] S. Tanase-Nicola and J. Kurchan, "Statistical-mechanical formulation of Lyapunov exponents", `cond-mat/0210380`. update

[2] Mason: expand