

In order to illuminate my problem, let us have a look at the 'prototype' 2-dimensional repeller, the single hyperbolic fixed point introduced in section 8.3, and try to calculate its escape rate:

$$f(z_1, z_2) = (\lambda_s z_1, \lambda_u z_2) \quad (1)$$

with  $0 < \lambda_s < 1$  and  $\lambda_u > 1$ . Of course the Perron-Frobenius operator

$$\begin{aligned} \mathcal{L}h(z_1, z_2) &= \int_{\mathbf{R}^2} dw_1 dw_2 \delta((z_1, z_2) - f(w_1, w_2)) h(w_1, w_2) \\ &= \frac{1}{\lambda_s \lambda_u} h(z_1/\lambda_s, z_2/\lambda_u) \end{aligned} \quad (2)$$

has smooth eigenfunctions  $\varphi_{n_1, n_2} = z_1^{-n_1-1} z_2^{n_2}$ , so I should be happy.

The drawback is that I must choose a finite sized neighbourhood of the repeller in order to define an escape rate. For the example given here, let this region be  $\mathcal{M} := [-1, 1] \times [-1, 1]$ , and for the escape rate I have to examine

$$\Gamma_n = \frac{1}{|\mathcal{M}|} \int_{\mathcal{M}} dw \int_{\mathcal{M}} dz \delta(z - f^n(w)), \quad (3)$$

and this is

$$\Gamma_n = \frac{1}{|\mathcal{M}|} \int_{\mathcal{M}} dz \tilde{\mathcal{L}}^n i(z). \quad (4)$$

with  $i(z) = 1$  and an operator

$$\tilde{\mathcal{L}}h(z_1, z_2) = \int_{\mathcal{M}} dw_1 dw_2 \delta((z_1, z_2) - f(w_1, w_2)) h(w_1, w_2) \quad (5)$$

which is definitely not the Perron-Frobenius operator defined in (2). The representation given by the second equality in (2) is valid for  $\tilde{\mathcal{L}}$  only if  $(z_1, z_2) \in f(\mathcal{M})$ , and  $\tilde{\mathcal{L}}h(z_1, z_2)$  is zero otherwise. The eigenfunctions of  $\mathcal{L}$  cannot be the eigenfunctions of  $\tilde{\mathcal{L}}$ . So what has the escape rate to do with the Perron-Frobenius operator? Why does it work?

Another question: Which linear combination of the  $\varphi_{n_1, n_2}$ 's yields  $i(z) = 1$ ? Smooth functions on  $\mathcal{M}$  do not belong to the somewhat strange function space spanned up by the  $\varphi_{n_1, n_2}$ 's.