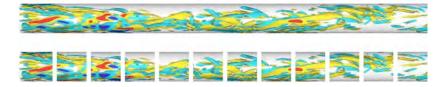
# kicked rotor cat in 1 spacetime dimension

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 $\begin{array}{l} \mbox{Georgia Tech} \\ \mbox{ChaosBook.org/overheads/spatiotemporal} \\ \rightarrow \mbox{Chaotic field theory slides} \end{array}$ 

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### building blocks od turbulence



#### have : a detailed theory of small turbulent cells

construct : the infinite state by coupling turbulent cells<sup>1</sup>

what would that theory look like ?

<sup>&</sup>lt;sup>1</sup>M. N. Gudorf et al., Spatiotemporal tiling of the Kuramoto-Sivashinsky flow, In preparation, 2021.

Field Theory should be Hamiltonian and energy conserving Quantum Mechanics requires it

because that is physics !

need a system as simple as the Bernoulli, but mechanical

so, we move on from running in circles,

to a mechanical rotor to kick.

#### next : a kicked rotor

Du mußt es dreimal sagen! — Mephistopheles

- what this is about
- 2 coin toss
- kicked rotor
- spatiotemporal cat
- bye bye, dynamics

field theory in 1 spacetime dimension

we now define

the cat map in 1 spacetime dimension

then we generalize to

d-dimensional spatiotemporal cat

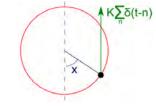
- cat map in Hamiltonian formulation
- cat map in Lagrangian formulation (so much more elegant!)

(1) the traditional cat

# time-evolution formulation

#### example of a "small domain" dynamics : a single kicked rotor

an electron circling an atom, subject to a discrete time sequence of angle-dependent kicks  $F(x_t)$ 



Taylor, Chirikov and Greene standard map

$$x_{t+1} - x_t = p_{t+1} \mod 1$$
  
 $p_{t+1} - p_t = F(x_t)$ 

 $\rightarrow$  chaos in Hamiltonian systems

#### the simplest example : a cat map evolving in time

force F(x) = Kx linear in the displacement x,  $K \in \mathbb{Z}$ 

 $x_{t+1} = x_t + p_{t+1} \mod 1$  $p_{t+1} = p_t + Kx_t \mod 1$ 

Continuous Automorphism of the Torus, or

### time-evolution cat map

a linear, area preserving map of a 2-torus onto itself

$$\begin{bmatrix} \phi_t \\ \phi_{t+1} \end{bmatrix} = J \begin{bmatrix} \phi_{t-1} \\ \phi_t \end{bmatrix} - \begin{bmatrix} 0 \\ m_t \end{bmatrix}, \qquad J = \begin{bmatrix} 0 & 1 \\ -1 & s \end{bmatrix}$$

for integer 'stretching'  $s = \operatorname{tr} J > 2$  the map is beloved by ergodicists : hyperbolic  $\Rightarrow$  perfect chaotic Hamiltonian dynamical system a cat is literally Hooke's wild, 'anti-harmonic' sister

## for *s* < 2 Hooke rules

local restoring oscillations around the sleepy z-z-z-zzz resting state

#### for s > 2 cats rule

exponential runaway wrapped global around a phase space torus

cat is to chaos what harmonic oscillator is to order

there is no more fundamental example of chaos in mechanics

(2) spatiotemporal cat

# lattice formulation

#### cat map in lattice formulation

replace momentum by velocity

$$p_{t+1} = (\phi_{t+1} - \phi_t)/\Delta t$$

obtain

$$\begin{bmatrix} \phi_t \\ \phi_{t+1} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & s \end{bmatrix} \begin{bmatrix} \phi_{t-1} \\ \phi_t \end{bmatrix} - \begin{bmatrix} 0 \\ m_t \end{bmatrix}$$

temporal lattice formulation is pretty<sup>2</sup> :

2-step difference equation

$$-\phi_{t+1} + s\phi_t - \phi_{t-1} = m_t$$

integer  $m_t$  ensures that

 $\phi_t$  lands in the unit interval

$$m_t \in \mathcal{A}, \quad \mathcal{A} = \{\text{finite alphabet}\}$$

<sup>2</sup>I. Percival and F. Vivaldi, Physica D 27, 373–386 (1987).

# spatiotemporal cat at every instant t, local in time

$$-\phi_{t+1} + \mathbf{s}\,\phi_t - \phi_{t-1} = \mathbf{m}_t$$

is enforced by the global equation

$$\mathcal{J}\Phi = \mathsf{M},$$

where

#### orbit Jacobian matrix

$$\mathcal{J} \Phi - \mathsf{M} = \mathsf{0}$$

with

$$\Phi = (\phi_{t+1}, \cdots, \phi_{t+n}), \quad \mathsf{M} = (m_{t+1}, \cdots, m_{t+n})$$

a lattice state, and a symbol block

and  $[n \times n]$  orbit Jacobian matrix  $\mathcal{J}$  is

$$-r + s \, \mathbb{1} - r^{-1} = \begin{pmatrix} s & -1 & & -1 \\ -1 & s & -1 & & \\ & -1 & & \ddots & \\ & & s & -1 \\ -1 & & -1 & s \end{pmatrix}$$

# think globally, act locally

solving the spatiotemporal cat equation

 $\mathcal{J}\Phi=\mathsf{M}\,,$ 

with the  $[n \times n]$  matrix  $\mathcal{J} = -r + s \mathbb{1} - r^{-1}$ can be viewed as a search for zeros of the function

$$F[\Phi] = \mathcal{J}\Phi - \mathsf{M} = 0$$

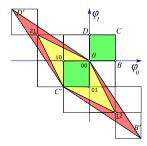
where the entire global lattice state  $\Phi_M$  is a single fixed point  $\Phi_M = (\phi_1, \phi_2, \cdots, \phi_n)$ 

in the *n*-dimensional unit hyper-cube



#### fundamental fact in action

temporal cat fundamental parallelepiped for period 2 square  $[0BCD] \Rightarrow \mathcal{J} =$  fundamental parallelepiped [0B'C'D']



$$N_2 = |\text{Det } \mathcal{J}| = 5$$

fundamental parallelepiped = 5 unit area quadrilaterals

a periodic point per each unit volume

#### spatiotemporal cat zeta function

is the generating function that counts orbits

substituting the Hill determinant count of periodic lattice states

 $N_n = \text{Det } \mathcal{J}$ 

into the topological zeta function

$$1/\zeta_{top}(z) = \exp\left(-\sum_{n=1} \frac{z^n}{n} N_n\right)$$

leads to the elegant explicit formula<sup>3</sup>

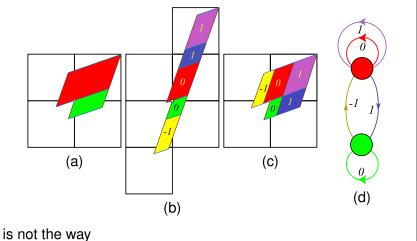
$$1/\zeta_{top}(z) = rac{1-sz+z^2}{1-2z+z^2}$$

# solved!

<sup>&</sup>lt;sup>3</sup>S. Isola, Europhys. Lett. **11**, 517–522 (1990).

# a side remark to experts

# slicing cats



Adler-Weiss generating partition of the unit torus is a distraction. Klein-Gordon is a deeper insight

### what continuum theory is temporal cat discretization of?

have

2-step difference equation

$$-\phi_{t+1} + \mathbf{s}\,\phi_t - \phi_{t-1} = m_t$$

discrete lattice

Laplacian in 1 dimension

$$\phi_{t+1} - \mathbf{2}\phi_t + \phi_{t-1} = \Box \phi_t$$

so temporal cat is an (anti)oscillator chain, known as

d = 1 Klein-Gordon (or damped Poisson) equation (!)

$$(-\Box + \mu^2) \phi_t = m_t, \qquad \mu^2 = s - 2$$

did you know that a cat map can be so cool?

a reminder slide, to skip : Helmholtz equation in continuum

# inhomogeneous Helmoltz equation

is an elliptical equation of form

$$(\Box + k^2) \phi(x) = -m(x), \qquad x \in \mathbb{R}^d$$

where  $\phi(x)$  is a  $C^2$  function, and m(x) is a function with compact support

for the  $\mu^2 = -k^2 > 0$  (imaginary *k*), the equation is known as the Klein-Gordon, Yukawa, or screened Poisson equation<sup>4</sup> equation

<sup>&</sup>lt;sup>4</sup>A. L. Fetter and J. D. Walecka, Theoretical Mechanics of Particles and Continua, (Dover, New York, 2003).

that's it! for spacetime of 1 dimension

lattice Klein-Gordon equation

$$(-\Box + \mu^2) \phi_t = m_t$$

solved completely and analytically!

# think globally, act locally - summary

the problem of determining all global solutions stripped to its bare essentials :

each solution a zero of the global fixed point condition

 $F[\Phi] = 0$ 

Output the orbit Jacobian matrix

$$\mathcal{J}_{ij} = \frac{\delta F[\Phi]_i}{\delta \phi_j}$$

fundamental fact

$$N_n = |\text{Det } \mathcal{J}| = \text{period-}n \text{ states}$$

 $\Rightarrow$  zeta function 1/ $\zeta_{top}(z)$ 

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